



# Structural and kinetic properties of *Rhodobacter sphaeroides* photosynthetic reaction centers containing exclusively Zn-coordinated bacteriochlorophyll as bacteriochlorin cofactors

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## ARTICLE INFO

### Article history:

Received 16 October 2013

Received in revised form 21 November 2013

Accepted 26 November 2013

Available online 5 December 2013

### Keywords:

Bacteriochlorophyll

Reaction center

Bacteriopheophytin

Electron transfer

*Rhodobacter sphaeroides*

Spectroscopy

## ABSTRACT

The Zn-BChl-containing reaction center (RC) produced in a *bchD* (magnesium chelatase) mutant of *Rhodobacter sphaeroides* assembles with six Zn-bacteriochlorophylls (Zn-BChls) in place of four Mg-containing bacteriochlorophylls (BChls) and two bacteriopheophytins (BPhe). This protein presents unique opportunities for studying biological electron transfer, as Zn-containing chlorins can exist in 4-, 5-, and (theoretically) 6-coordinate states within the RC. In this paper, the electron transfer perturbations attributed exclusively to coordination state effects are separated from those attributed to the presence, absence, or type of metal in the bacteriochlorin at the H<sub>A</sub> pocket of the RC. The presence of a 4-coordinate Zn<sup>2+</sup> ion in the H<sub>A</sub> bacteriochlorin instead of BPhe results in a small decrease in the rates of the P<sup>+</sup> → P<sup>+</sup>H<sub>A</sub><sup>−</sup> → P<sup>+</sup>Q<sub>A</sub><sup>−</sup> electron transfer, and the charge separation yield is not greatly perturbed; however coordination of the Zn<sup>2+</sup> by a fifth ligand provided by a histidine residue results in a larger rate decrease and yield loss. We also report the first crystal structure of a Zn-BChl-containing RC, confirming that the H<sub>A</sub> Zn-BChl was either 4- or 5-coordinate in the two types of Zn-BChl-containing RCs studied here. Interestingly, a large degree of disorder, in combination with a relatively weak anomalous difference electron density was found in the H<sub>B</sub> pocket. These data, in combination with spectroscopic results, indicate partial occupancy of this binding pocket. These findings provide insights into the use of BPhe as the bacteriochlorin pigment of choice at H<sub>A</sub> in both BChl- and Zn-BChl-containing RCs found in nature.

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## 1. Introduction

Photosynthetic organisms utilize light energy to generate ATP via photophosphorylation. In order to trap light energy, reaction center (RC) proteins convert electronic excitation energy from photon absorption into charge separation. This means of charge separation relies on a cascade of ultrafast electron transfer reactions that proceeds through a series of cofactors buried within the protein. The photosynthetic RC of the bacterium *Rhodobacter sphaeroides* is the most extensively characterized and best understood RC in terms of the mechanisms of electron transfer events that result in charge separation [1–3]. The crystal structure of this protein revealed that the cofactors are arranged into two structurally largely symmetrical branches (A and B) housed within the three polypeptide subunits, designated L, M and H [4,5]. Although the

two pigment branches are structurally symmetrical, they are functionally asymmetrical because the electron transfer events occur almost exclusively along the A branch, with the B branch playing a role only in the final electron transfer step.

In the initial stages of charge separation, a dimer of bacteriochlorophyll (BChl) molecules (P<sub>A,B</sub>) located at the symmetry axis of the two pigment branches enters an excited state as a result of direct light absorption or, more commonly, excitation energy transfer from the surrounding antenna complex LH1. This excited state (P<sup>\*</sup>) decays via a 3 ps electron transfer to a nearby BChl on the A branch in a protein pocket, termed B<sub>A</sub>, generating a P<sup>+</sup>B<sub>A</sub><sup>−</sup> state within the RC. Subsequently, the B<sub>A</sub><sup>−</sup> radical anion (radical anions in the RC are subsequently referred to simply as anions) reduces a nearby bacteriopheophytin (BPhe) molecule in the H<sub>A</sub> pocket within a single picosecond, generating a P<sup>+</sup>H<sub>A</sub><sup>−</sup> state. The P<sup>+</sup>H<sub>A</sub><sup>−</sup> state is relatively long-lived compared to P<sup>\*</sup> and P<sup>+</sup>B<sub>A</sub><sup>−</sup>, lasting ~200 ps before an electron transfer reaction from the H<sub>A</sub><sup>−</sup> anion to a nearby quinone, Q<sub>A</sub>. Once the RC is poised in the P<sup>+</sup>Q<sub>A</sub><sup>−</sup> state, a final electron transfer reaction reduces the quinone on the B branch, Q<sub>B</sub>. At this point, the oxidized dimeric BChls must receive an electron from a c-type cytochrome in the periplasm before the RC is set for another turnover, which results

**Abbreviations:** E, *Escherichia*; R, *Rhodobacter*; A, *Acidiphilium*; RC, reaction center; BChl, bacteriochlorophyll; Zn-BChl, [Zn]-bacteriochlorophyll; BPhe, bacteriopheophytin; LDAO, lauryldimethylamine-*n*-oxide; EADS, evolution associated difference spectra; LH1, light harvesting 1; LH2, light harvesting 2

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in the double reduction of  $Q_B$ . These electron transfers are coupled to protonation from the cytoplasm, converting  $Q_B$  from a quinone to a quinol. The quinol leaves the RC and migrates within the membrane to the cytochrome  $bc_1$  complex, which ultimately generates a proton gradient that powers an  $F_0F_1$  ATP synthase [6].

Spectroscopic advances over the past few decades have allowed access to the femtosecond time domain, facilitating studies of electron transfer in the RC [7]. Such experiments have shed light on some of the factors that control the rates of electron transfer, such as cofactor composition and cofactor midpoint potential, as well as the surrounding protein environment. It is generally accepted that the driving force for efficient electron transfer from  $B_A$  to  $H_A$  is dominated by the energy difference ( $\Delta G$ ) between the two cofactors, resulting from their different structures (BChl vs BPhe). This is true for all type II reaction centers [8]. However, a RC assembled with Zn-BChl in place of all BChls and BPhe was shown to have nearly wild type (WT) electron transfer rates [9,10]. The electron transfer rates in this Zn-RC were surprising because of the similarity of this RC's A-branch cofactor composition to the so-called  $\beta$  mutant RC originally characterized by Kirmaier et al. [11]. In the  $\beta$  mutant RC, an engineered mutation of leucine to histidine at position 214 of the M subunit ((M) L214H) resulted in a molecule of BChl being present in the  $H_A$  pocket, and this RC exhibited a ~60% yield of charge separation compared to the WT RC within a 2 ns interval after excitation [11].

In both the Zn-RC and  $\beta$  mutant (Mg-BChl) RC, electrons proceed from dimeric (Zn-) BChls to a monomeric (Zn-) BChl at  $B_A$  and onward to yet another (Zn-) BChl at  $H_A$ . Apparently the  $H_A$  Zn-BChl is not converted to BPhe because the  $Zn^{2+}$  ion is more tightly bound to the bacteriochlorin macrocycle than a  $Mg^{2+}$  [12]. Indeed, Zn-BChl is found in the RC and antenna complexes of an acidophilic photosynthetic bacterium, *Acidiphilium rubrum* [13], which lives at low pH (2.5–6) where Mg-containing BChl would potentially be converted to BPhe [13]. Lin et al. postulated that in addition to the difference of a  $Mg^{2+}$  vs.  $Zn^{2+}$  ion in the  $H_A$  bacteriochlorin, the difference in coordination state of the metal in the bacteriochlorin in the  $H_A$  pocket was responsible for a change in electron transfer rates by a factor of about two in these two RCs that contain a Zn- or Mg-BChl in place of BPhe in the  $H_A$  pocket [9].

In the case of the  $\beta$  mutant RC studied previously, the central  $Mg^{2+}$  in the  $H_A$  BChl was thought to be 5-coordinate [11]. In the Zn-RC, the analogous  $H_A$  cofactor was proposed to be 4-coordinate, consistent with an absorption spectrum more similar to BPhe than a 5-coordinate BChl in the  $Q_X$  transition region [9,14]. Previous work has indicated a correlation between  $Q_X$  transition energy and redox potential in metallo-bacteriochlorins, implying that the redox potential of the Zn-BChl in the  $H_A$  pocket may be close to that of BPhe [15,16].

In a hole-burning study of the Zn- and Zn- $\beta$ -RC (the Zn- $\beta$ -RC contains the (M) L214H mutation of the  $\beta$  mutant engineered in a *bchD* mutant background), it was proposed that the quantum yield of  $P^+Q_A^-$  formation decreased by 40% in the Zn- $\beta$ -RC, relative to the Zn-RC. This estimate was based on the assumption that the  $P^+Q_A^-$  lifetime is similar in the two RCs [10]. The  $H_A$  Zn-BChl was proposed to be 5-coordinate in the Zn- $\beta$ -RC, based on the presence of the (M) L214H side chain and the disappearance of a 542 nm peak (attributed to the  $H_{A,B}$   $Q_X$  transition) in low temperature absorption spectra. However, this was a larger absorbance change than expected, as it was thought that the  $H_{A,B}$   $Q_X$  transition peak should decrease by only ~50% as a result of  $H_A$  Zn-BChl penta-coordination [10]. Several explanations were offered, because there were uncertainties about whether the  $H_B$  or carotenoid cofactor was present, and whether the  $H_B$  Zn-BChl was present but as 5-coordinate (perhaps with water as the fifth  $Zn^{2+}$  ligand) [10]. Furthermore, the forward and back electron transfer rates in the pathway between P and  $Q_A$  were not addressed.

Here, the effects of coordination state of a metal-containing bacteriochlorin in the  $H_A$  pocket on RC structure and function are explored in more detail by directly comparing two RCs containing six Zn-BChls, and that differ only in whether leucine or histidine is present at M214.

The Zn-BChl-containing RC with (M)214-leucine (Zn-RC) is compared to the Zn-BChl-containing RC with M214-histidine (Zn- $\beta$ -RC) using absorption spectroscopy, X-ray anomalous scattering and structure analysis of protein crystals, and femtosecond transient absorption spectroscopy. For clarity, the RCs in this study are henceforth called the WT-RC, WT- $\beta$ -RC (meaning the (M) L214H mutation), Zn-RC, and Zn- $\beta$ -RC. Using these RCs, the effects attributed to the presence or absence of a central metal at  $H_A$  were separated from those attributed to protein ligation and hence coordination state. Our results suggest that there is a small decrease in charge separation yield due to the mere presence of a  $Zn^{2+}$  ion in the bacteriochlorin at  $H_A$ , and a larger yield loss upon formation of 5-coordinate Zn-BChl. Furthermore, we report structural differences at  $H_B$  that appear to result from the substitution of Zn-BChl for BChl. These data are discussed in the context of (i) the consequences of metal coordination vs. axial ligation at  $H_A$  on RC function, (ii) the preference for BPhe in photosynthetic RCs, even in an organism that naturally incorporates Zn-BChl into the RC, and (iii) the general effects of alternative pigment production in photosynthetic bacteria.

## 2. Materials and methods

### 2.1. Bacterial strains, plasmids and cultures

Starter cultures of *R. sphaeroides* were grown in RCV medium aerobically in the dark [17]. These cultures were used to seed small chemotrophic RLB medium [18] cultures which were grown semi-aerobically in the dark at 150 RPM in 250 mL conical flasks. These cultures were then used to seed 5.6 L of RLB media (25 mL inocula added to  $4 \times 2$  L conical flasks each containing 1.4 L of media), which were shaken at 150 RPM in the dark. *Escherichia* (*E.*) *coli* strain DH10B was used for cloning and site-directed mutagenesis. *R. sphaeroides* strains  $\Delta RCLH$  and  $\Delta RCLH/\Delta bchD$  were previously described [10,19]. The previously described plasmids pRS1 and pRS1:::(M) L214H were transferred from *E. coli* strain DH10B to *R. sphaeroides* strain  $\Delta RCLH/\Delta bchD$  by conjugation [20].

### 2.2. Isolation of reaction centers

His-tagged RCs were isolated from 5.6 L of culture according to the procedure of Saer et al. [20], essentially a modified version of the method of Goldsmith and Boxer [20,21]. RCs were dialyzed against 10 mM Tris, 0.025% LDAO and stored at  $-80^\circ\text{C}$  for later use. To obtain sufficient amounts of Zn-BChl-containing RCs, the cell culture volume was scaled up to 72 L because of the low yield of these RCs from *bchD* mutant strains. RCs used in protein crystallography were never frozen and subsequently purified by FPLC using a SourceQ column prior to hanging drop experiments.

### 2.3. Ground state absorption spectroscopy

Ground state absorption spectra of purified RCs in 10 mM Tris (pH 8.0), 0.025% LDAO mixed with an equal volume of glycerol were obtained at 10 K with an ARS helium cryostat and a Cary 6000i spectrophotometer using the same methods and instrumentation described by Saer et al. [20]. Where necessary, a monotonic increase in background absorbance (500 to 730 nm) resulting from light-scattering was corrected by subtracting a linear function.

### 2.4. Time-resolved spectroscopy

The femtosecond transient absorption spectrometer utilized was the same setup described in Pan et al. [22]. Transient absorption changes from a spectral range of 500–980 nm, and a time range of 0.5 ps before and 6 ns after the excitation pulses, were recorded. RC samples were loaded in a 9 cm diameter spinning wheel (~10 revolutions per second) with an optical path length of 1.2 mm at an optical density of ~0.8 at

800 nm. Global analysis of the datasets was performed with ASUFIT, a locally-written program created in the MATLAB environment.

### 2.5. X-ray crystallography

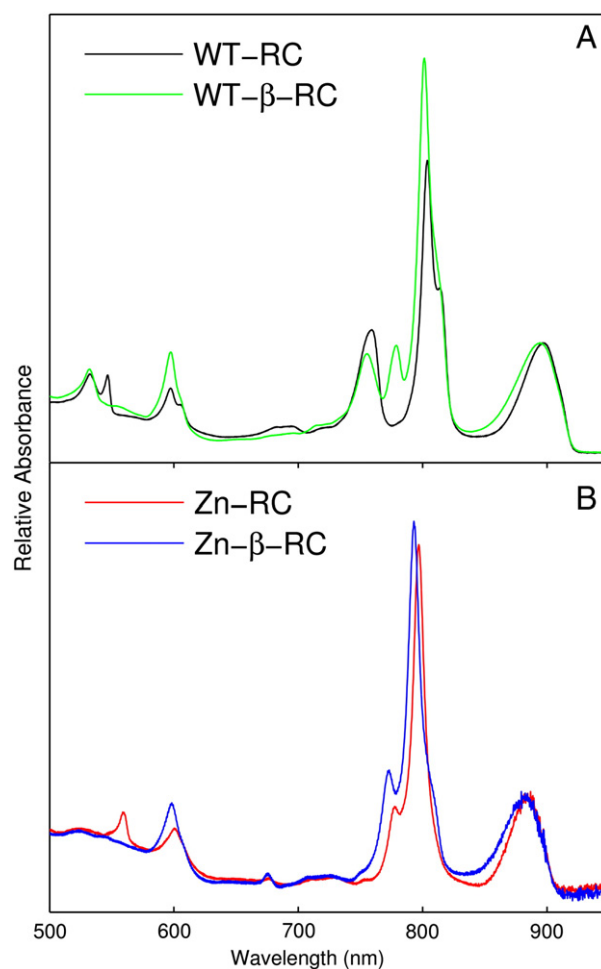
Crystals of the Zn-RC were formed using the hanging drop technique outlined in Saer et al. [20]. Diffraction data were collected at Stanford Synchrotron Light Source on beamline 7.1 using a wavelength of 1.12709 Å and processed with Mosflm and Aimless to 2.85 Å resolution [23,24]. Anomalous differences were detected to 6 Å resolution. All data sets were isomorphous with wild type RC (PDB entry 2J8C), which was used as the starting point for refinement with removal of the  $Mg^{2+}$  ions and for Zn-β-RC the side chain atoms of (M)214. Initial phases were obtained by limited refinement with Refmac5 [25] and  $F_o - F_c$  and anomalous difference maps were computed with Coot [26] and FFT, respectively. Based on the peaks observed in both maps, Zn-BChl was modeled at each of the P, B, and H cofactor sites in both the A and B branches. Density for a histidine side chain was observed at residue (M)214 in the Zn-β-RC and modeled accordingly using the program Coot. Refinement of the structures was continued with Refmac5 and the removal of solvent atoms with unrealistic B-factors.

## 3. Results

### 3.1. Low temperature absorption spectroscopy

The ground state absorption spectra of the WT-RC, WT-β-RC, Zn-RC and Zn-β-RC at 10 K are shown in Fig. 1. The spectra were normalized to the amplitude of the  $Q_y$  band of the dimeric BChls ( $P_{A,B}$ ). Characteristic absorption differences between the WT-RC and the WT-β-RC at cryogenic temperatures have been described [27]. The important features that distinguish these two RCs are the loss of the 542 nm  $Q_x$  absorption band and diminished amplitude of the 765 nm  $Q_y$  band in the WT-β-RC. These changes are accompanied by an increase in amplitude of the 600 nm  $Q_x$  band, and the emergence of a 785 nm  $Q_y$  band. Collectively, these absorption changes indicate replacement of the  $H_A$  BPhe in the WT-β-RC with a molecule of BChl. The  $Q_y$  transition differences in absorption between the Zn-β-RC and the Zn-RC are more subtle than the differences between the WT-β-RC and the WT-RC. The  $P_{A,B}$  band in the Zn-β-RC is slightly broadened relative to the Zn-RC, and the  $B_{A,B}$  and  $H_{A,B}$   $Q_y$  bands of the Zn-β-RC appear to be shifted ~2 nm to the blue. In the  $Q_x$  transition region of the Zn-β-RC spectrum, several features indicate the conversion of the 4-coordinate Zn-BChl at  $H_A$  in the Zn-RC to 5-coordinate in the Zn-β-RC: the increased amplitude of the ~590 nm peak corresponding to a 5-coordinate Zn-BChl, accompanied by the disappearance of a peak at ~560 nm, where the 4-coordinate Zn-BChl absorbs in the Zn-RC [9,10]. The  $H_B$  Zn-BChl content in the Zn-RC appears to be low, because the peak at ~560 nm seems too small to represent contributions from two Zn-BChls (compare this region in the Zn-RC to the WT-RC, which shows two distinct peaks of roughly the same size).

The bandwidth of  $P_{A,B}$  absorption in the  $Q_y$  transition for both the WT-β and Zn-β-RC appears slightly broadened on the shorter wavelength side relative to their non-β counterparts. Therefore, although (M)H214 should not be in the immediate vicinity of the macrocycle component of the dimeric (Zn-)BChls, it appears that this change slightly perturbs this region. The highest resolution crystal structure available for the RC shows that (M) L214 lies close to the phytol tail of the  $P_A$  BChl [28]. Thus, it appears that this bandwidth increase may be a consequence of an interaction between the histidine residue at (M)214 and the  $P_A$  tail. Alternatively, the (M)L214H mutation may induce a small perturbation in the membrane-spanning helix (helix D) in which this residue resides. Because this helix also contains residue (M)H202, which acts as the fifth coordinate to  $P_B$  [29], changes in (M)214 may contribute to changes in  $P_{A,B}$   $Q_y$  absorbance through minor variations of the  $P_B$ -(M)H202 bond.



**Fig. 1.** Ground state absorption spectra obtained at 10 K for (A) the WT (black), and β mutant RCs (green), and (B) the Zn-RC (red) and Zn-β RCs (blue). The spectra were normalized to the P865 absorption maximum. The spectrum of the Zn-RC was corrected for scatter by subtracting a straight line from 500 to 730 nm.

### 3.2. X-ray crystallography and anomalous scattering

To elucidate the type, coordination state and occupancy of the bacteriochlorin cofactors in Zn-BChl-containing RCs, the crystal structures of the Zn-RC and Zn-β-RC were solved to a resolution of 2.85 Å. The refinement statistics are listed in Table 1. Overall, the structures resemble that of the native RC (PDB entry 2J8C). In both the Zn-RC and Zn-β-RC structures, the iron, carotenoid and all the bacteriochlorin cofactors are present.

An electron density map was computed from the anomalous difference data collected above the Zn-edge of the two RC structures to reveal the location of  $Zn^{2+}$  ions within the two structures (Table 2). A peak greater than  $4\sigma$  was observed in the anomalous maps at the center of all bacteriochlorin-type cofactors indicative of the presence of  $Zn^{2+}$  ions at these sites. At the wavelength used for data collection, the anomalous signal from  $Mg^{2+}$  is weak and would not significantly contribute to the density observed in the map. The signal from iron can be used as an internal standard and indeed peaks at  $8.7\sigma$  and  $10\sigma$  were observed at the iron site in the Zn-RC and Zn-β-RC structures, respectively. In the Zn-RC structure, five of the bacteriochlorin cofactors appear to be bound at full occupancy based on refined  $Zn^{2+}$  ion B-factors of  $58 \text{ Å}^2$  or less and anomalous map peaks  $> 12\sigma$ . In contrast, the larger B-factor of the Zn atom at  $H_B$  site ( $86 \text{ Å}^2$ ), in combination with the lower anomalous map peak ( $7\sigma$ ) imply a lower occupancy and/or a higher degree of disorder of this cofactor. The  $H_B$  cofactor in the Zn-β-RC structure exhibits even lower occupancy and/or greater disorder (Table 2). These

**Table 1**  
X-ray data collection and refinement statistics for Zn-RC and Zn-β-RC.

Structure		Zn-RC	Zn-β-RC
Accession code		4N7K	4N7L
Unit cell parameters	<i>a</i> = <i>b</i>	139.45 Å	139.59 Å
	<i>c</i>	184.18 Å	184.05 Å
	Resolution (Å)	73.23–2.85 (2.90–2.85)	69.79–2.85 (2.90–2.85)
<i>R</i> <sub>merge</sub>		0.105 (0.972)	0.098 (0.705)
<i>I</i> / <i>σI</i>		10.5 (1.5)	12.6 (2.4)
Multiplicity		6.2 (4.7)	7.0 (6.2)
Completeness (%)		96 (95)	87.4 (85.6)
<i>R</i> <sub>work</sub>		0.182	0.167
<i>R</i> <sub>free</sub>		0.242	0.216
Wilson B-factor (Å <sup>2</sup> )		53.1	47.0
Overall B-factor (Å <sup>2</sup> )		58.2	54.8
Coordinate error <sup>a</sup> (Å)		0.21	0.17
Bond length rmsd (Å)		0.011	0.013
Bond angle rmsd (°)		2.6	2.7

Numbers in parentheses reflect statistics for the highest resolution shell.

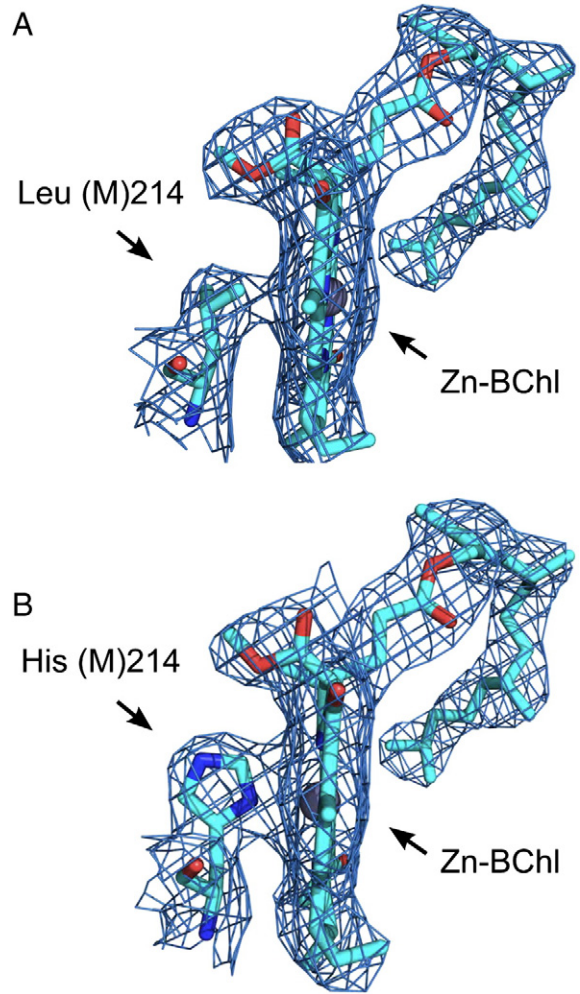
<sup>a</sup> Coordinate error is the estimated standard uncertainty from maximum likelihood refinement.

values for the H<sub>B</sub> cofactor in both the Zn-RC and Zn-β-RC, relative to the values for the other bacteriochlorin cofactors, indicate that the H<sub>B</sub> pocket is not well-occupied, and perhaps less well-occupied in the Zn-β-RC than in the Zn-RC. Such a decrease in occupancy may provide a rationale for the absence of a Q<sub>x</sub> transition H<sub>B</sub> peak (Fig. 1 and Neupane et al.) [10]. Interestingly, the B-factors of the amino acid residues surrounding the H<sub>B</sub> cofactor were similar to those in the rest of the protein, indicating that the observed disorder was centered on the Zn-Bchl molecule itself, and not the surrounding protein.

In comparison to the H<sub>B</sub> site, the B<sub>B</sub>, P<sub>A/B</sub>, and A-branch Zn-BChls of the Zn-β-RC and Zn-RC appear to be relatively well-ordered and fully occupied, based on the anomalous signal intensity and B-factor values, which are similar to the mean values of each respective structure (Table 2). A side by side comparison of the H<sub>A</sub> cofactors for the Zn-RC and Zn-β-RC can be seen in Fig. 2. For Zn-Bchl in sites with a His residue available for coordination, the imidazole ring is observed 2.1 to 2.4 Å from the Zn<sup>2+</sup> ion. At the Zn-RC H<sub>A</sub> site (with a Leu present), the side chain is separated by 3.2 Å from the Zn<sup>2+</sup> ion of the Zn-Bchl. Although at 2.85 Å resolution for these structures the error estimate of these distances is ±0.3 Å, at all of the Zn-BChls the electron density maps indicate the absence of an externally-derived metal ligand, such as water. Therefore, we suggest that the major structural difference between the A-branch electron transfer carriers in the Zn-RC and Zn-β-RC is the coordination state of the Zn<sup>2+</sup> ion in the H<sub>A</sub> Zn-Bchl, which is 4-coordinate in the Zn-RC and 5-coordinate in the Zn-β-RC.

### 3.3. Time-resolved spectroscopy

Upon RC excitation with an 860 nm pulse, P forms an excited state (P\*) resulting in stimulated emission which dominates the absorbance change signal near 925 nm (Fig. 3). P\* then decays by charge separation on the picosecond time scale forming the P<sup>+</sup>H<sub>A</sub><sup>−</sup> state, or possibly a mix of P<sup>+</sup>B<sub>A</sub><sup>−</sup> and P<sup>+</sup>H<sub>A</sub><sup>−</sup>, also called P<sup>+</sup>I<sub>A</sub><sup>−</sup> [27]. Global analysis of the time-

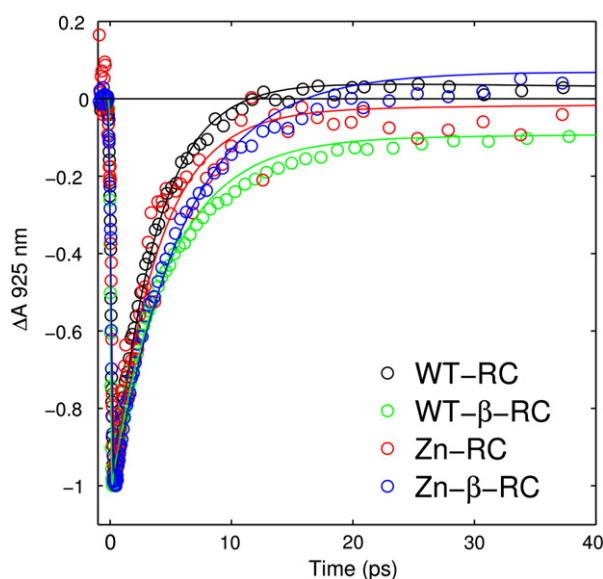


**Fig. 2.** Stick models and electron density (mesh) of the H<sub>A</sub> cofactor and axial (M)214 residues from the crystal structures of the (A) Zn-RC and (B) Zn-β-RC. The distances from the axial residues to the macrocycle centers are 3.9 Å and 2.1 to 2.4 Å for the Zn-RC and Zn-β-RC, respectively. Color code: teal, carbon atoms; blue, nitrogen atoms; red, oxygen atoms.

resolved spectra for all four RCs was performed over the Q<sub>x</sub> wavelength region from 500 to 750 nm and Q<sub>y</sub> region from 700 to 960 nm, by fitting to a sequential reaction model (A → B → C ...). The time constants of each reaction step, in addition to the yields of charge separation and calculated ground state recombination rates, are compiled in Table 3. The fit in the Q<sub>y</sub> region for the decay of P\*, representative of forward electron transfer to the nearby B<sub>A</sub> cofactor (and subsequently to H<sub>A</sub>), yielded time constants of 3.8 ps for the WT-RC, 3.9 ps for the Zn-RC, 4.8 ps for the WT-β-RC, and 5.5 ps for Zn-β-RC (Table 3). Thus, initial electron transfer takes place on the picosecond time scale for all four RCs

**Table 2**  
Anomalous map statistics for metals in the Zn-RC and Zn-β-RC X-ray diffraction datasets.

Site	Zn-RC				Zn-β-RC			
	OmitFo-Fc peak (σ)	Anomalous map peak (σ)	Zn B-factor (Å <sup>2</sup> )	Leu/His-Zn distance (Å)	Omit Fo-Fc peak (σ)	Anomalous map peak (σ)	Zn B-factor (Å <sup>2</sup> )	Leu/His-Zn distance (Å)
P <sub>A</sub>	25	13	50	2.1	25	16	43	2.2
P <sub>B</sub>	25	13	50	2.1	24	14	44	2.2
B <sub>A</sub>	25	14	44	2.2	25	15	39	2.3
B <sub>B</sub>	24	14	39	2.4	25	16	37	2.5
H <sub>A</sub>	20	12	58	3.2	26	15	43	2.3
H <sub>B</sub>	13	7	86	3.6	11	4	113	3.9
Fe	24	9	45	–	24	10	41	–



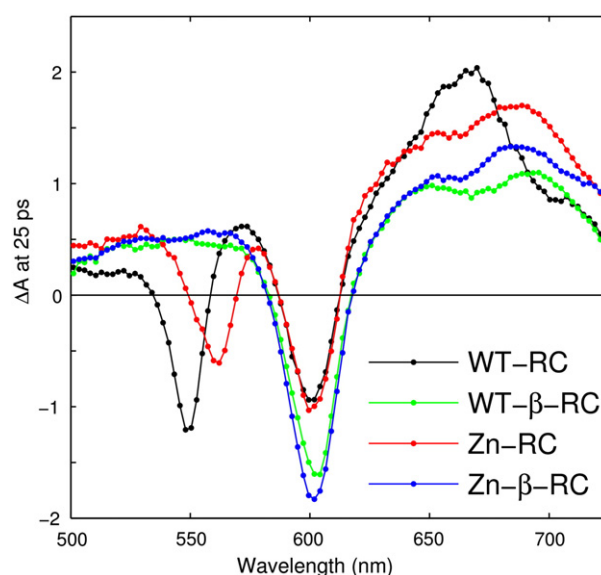
**Fig. 3.** Kinetic traces of stimulated emission decay for WT,  $\beta$  mutant, Zn-, and Zn- $\beta$  RCs at 925 nm. Data points are shown as open circles, whereas the fits are shown as solid lines. The data were normalized to the maximum level of bleaching.

although the rates of the WT-RC and Zn-RC group together, separately from the rates of the WT- $\beta$ -RC and Zn- $\beta$ -RC.

Previous absorption and fluorescence studies have shown that it is possible to resolve a second decay phase from the  $P^*$  kinetics for datasets in which a high signal-to-noise ratio of the data is achieved [30,31]. Our fitting revealed the presence of a second decay phase in all four RCs studied. While the amplitudes of this phase were small compared to the initial phase of  $P^*$  decay (23% for the WT-RC; 31% for the WT- $\beta$ -RC; 34% for the Zn-RC; and 32% for the Zn- $\beta$ -RC), the presence of this component was required to obtain a good fit of the data. The lifetimes of the longer decay phase were  $11 \pm 5$  ps for the WT-RC,  $11 \pm 7$  ps for the Zn-RC,  $18 \pm 5$  ps for the WT- $\beta$ -RC, and  $22 \pm 5$  ps for the Zn- $\beta$ -RC (Table 3).

Transient absorption difference spectra from 500 to 725 nm for all RCs at 25 ps after excitation are shown in Fig. 4. The spectra were normalized to the  $P^*$  bleaching maximum near 600 nm at 0.5 ps. In the Zn-RC, there is an absorption decrease near 560 nm, attributed to bleaching of the Zn-Bchl in the  $H_A$  site because this wavelength corresponds to the  $Q_x$  absorption maximum of this cofactor in the ground state (Fig. 1). Like the ground state transition, the magnitude of this bleaching is substantially less than that of the corresponding  $H_A$  bleaching in the WT-RC at ~545 nm (we assume similar extinction coefficients). In contrast, the magnitude of absorbance decrease near 600 nm, which represents ground state bleaching of P on this time scale, is nearly identical in the WT-RC and Zn-RC.

In the Zn- $\beta$ -RC there is no ground state absorbance in the 545 to 560 nm region (Fig. 1), and therefore no ground state bleaching in this region of the difference spectrum (Fig. 4). However the bleaching at ~600 nm is increased compared to both the WT-RC and Zn-RC. This



**Fig. 4.** Transient absorption difference spectra of the  $Q_x$  transitions of WT and mutant RCs 25 ps after excitation, representative of the  $P^+I_A^-$  state. The data were normalized to the  $Q_x$  bleaching maximum of  $P^*$  (~600 nm) at 0.5 ps post-excitation.

is consistent with ground state absorption measurements of the Zn- $\beta$ -RC at low temperatures, which show that the  $H_A$  Zn-Bchl absorbs maximally at 598 nm, overlapping the ground state bands of P,  $B_A$  and  $B_B$  (Fig. 1) [10]. The red shift from ~545 to ~600 nm in both the ground state absorption spectrum and the difference spectrum of the WT- $\beta$ -RC results from conversion of BPhe to His-coordinated Bchl, whereas in the Zn- $\beta$ -RC the analogous shift is consistent with a Zn-Bchl coordination state change at  $H_A$  from 4- to 5-coordinate.

A prominent feature of the  $P^+I_A^-$  absorption spectrum of all these RCs is the large broadband absorption increase in the ~620 to 725 nm region (Fig. 4), primarily due to anion absorption in this region. The peak of this absorbance increase shifts from ~660 nm in the WT-RC to ~680 nm in the WT- $\beta$ -RC, as shown previously [27]. In both the Zn-RC and Zn- $\beta$ -RC, the portion of this maximum is intermediate in wavelength between that of the WT-RC and the WT- $\beta$ -RC. The fact that the two Zn-Bchl-containing RCs share a  $P^+I_A^-$  absorption maximum at a similar wavelength, and which differs from that of both the WT-RC and the WT- $\beta$ -RC indicates that the position of the anion absorption peak is sensitive to both the presence and identity of the metal associated with the  $H_A$  bacteriochlorin, but insensitive to its coordination state.

The absorbance between ~620 and 725 nm associated with the  $I_A^-$  anion decays in the WT-RC with a time constant of 200 to 250 ps, as shown for the 675 nm wavelength in Fig. 5, in agreement with previous work [9,30,32,33]. The decay of the  $I_A^-$  anion band in the Zn-RC has a time constant of ~290 ps (Fig. 5, Table 3), as seen previously [9]. In both the WT- $\beta$ -RC mutant and Zn- $\beta$  RCs, much longer time constants of ~560 ps and ~640 ps, respectively, were resolved from the fitting (Table 3). Note that minima of larger uncertainty ( $\pm 50$  ps) were found in the fits for the WT- $\beta$ -RC and Zn- $\beta$ -RC compared to the WT

**Table 3**

Calculated transient state lifetimes, reaction rates and final yields of charge separation for RCs in this study based on global analysis.

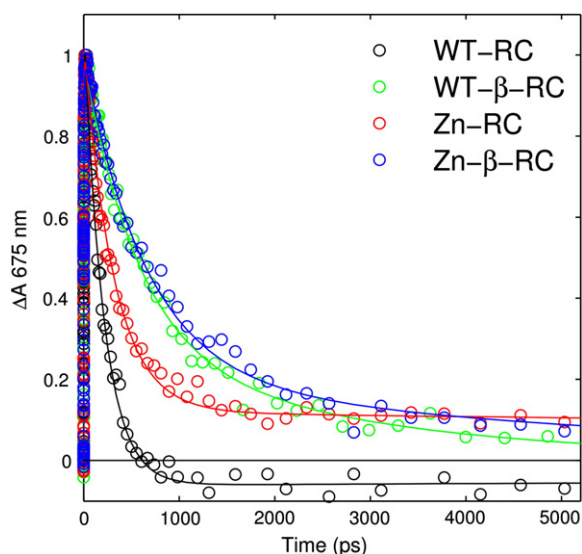
Sample	$P^*$ (fast)	$P^*$ (slow)	$P^+I_A^-$	$\tau_{IQ}^a$	$\tau_{rec}^b$	$\phi^c$ ( $P^+Q_A^-$ )
WT-RC	$3.8 \pm 0.2$ ps	$11 \pm 5$ ps	$210 \pm 30$ ps	0.2 ns	$20 \text{ ns}^d$	100%
WT- $\beta$ -RC	$4.8 \pm 0.5$ ps	$18 \pm 5$ ps	$560 \pm 100$ ps	$4.0 \pm 0.7$ ns	$0.65 \pm 0.1$ ns	14%
Zn-RC	$3.9 \pm 0.7$ ps	$11 \pm 7$ ps	$290 \pm 50$ ps	$0.4 \pm 0.1$ ns	$1.1 \pm 0.2$ ns	74%
Zn- $\beta$ -RC	$5.5 \pm 0.2$ ps	$22 \pm 5$ ps	$640 \pm 90$ ps	$6.4 \pm 0.9$ ns	$0.71 \pm 0.1$ ns	10%

<sup>a</sup> Calculated rate of forward electron transfer from  $I_A^-$  to  $Q_A$  using Eq. (2).

<sup>b</sup> Calculated rate of recombination using Eq. (1).

<sup>c</sup> Calculated yield of  $P^+Q_A^-$  based on the amount of  $P^+$  bleaching amplitude (measured at 860 nm) at the end of the measurement window compared to 45 ps post excitation.

<sup>d</sup> Based on measurements in which  $Q_A$  was removed from the RC.

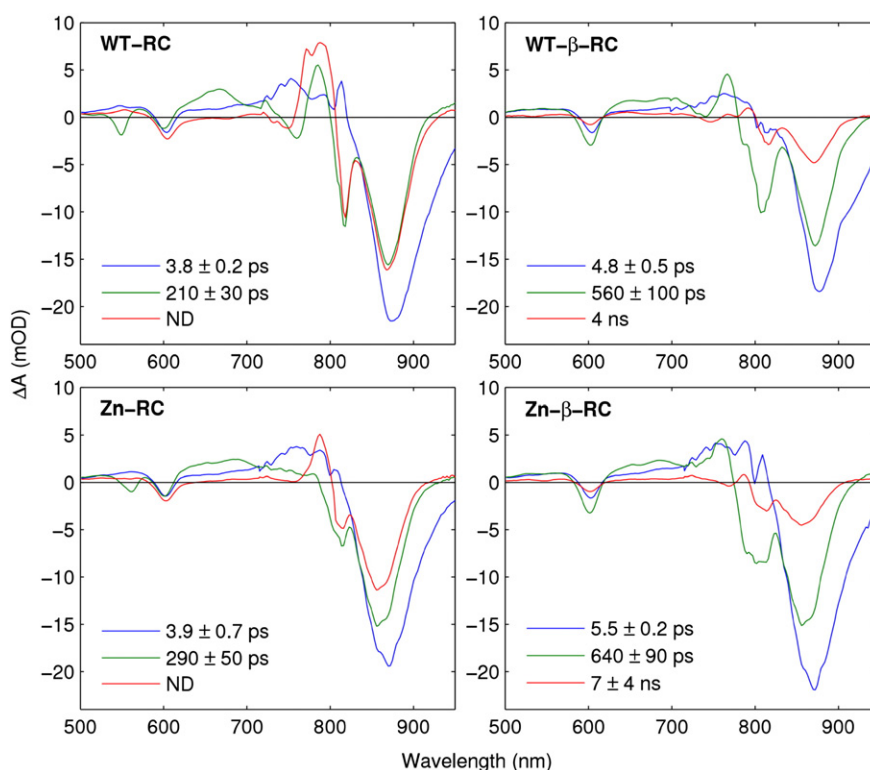


**Fig. 5.** Kinetics of  $P^+I_A^-$  state decay at 675 nm for WT and mutant RCs measured up to 5500 ps. Spectra were normalized to the maximum amplitude increase at 675 nm. Open circles represent raw data points, and solid lines are representative of fits to the data.

and Zn-RCs. The analysis of this complication will be given in the Discussion section.

In Fig. 6, the fitting results for each type of RC are represented as evolution-associated difference spectra (EADS) [34]. This representation allows simplified visualization of the spectral progression associated with the decay times of the fit, and is not meant to depict the spectra of the actual excited and charge separated state populations

(which would require the application of a specific kinetic model). As can be seen, the values of the time constants for the hundreds of picoseconds EADS for the WT-RC and Zn-RC (210 ps and 290 ps, respectively) are more similar to each other than to the time constants found for the WT-β-RC and the Zn-β-RC (560 ps and 640 ps, respectively). Furthermore, the EADS with the longest time constant in both the WT-RC and Zn-RC was non-decaying (*i.e.*, greater than 3 times the time window of measurement), whereas the WT-β-RC and the Zn-β-RC did not contain such a long-lived component. Thus, the kinetic behavior of the 4-coordinate Zn-BChl at  $H_A$  is more like that of a BPhe than a 5-coordinate (Zn-)BChl. All of the mutants, including the Zn-RC, however, appear to have lower  $P^+Q_A^-$  yields than the WT-RC on the nanosecond time scale. This can be seen in Fig. 6 by comparing the bleaching of P near 860 nm in the green EADS component that has a lifetime of hundreds of picoseconds, with that of the red (nanosecond, or non-decaying) component. In the WT-RC, there is no significant decrease in the amplitude of the bleaching signal over this time, as would be expected for 100% yield of charge separation. In contrast, each of the mutants shows a partial decrease in bleaching amplitude. In the case of the Zn-RC, this decrease is relatively small (74% of the bleaching amplitude remains; Table 3), however this amplitude is still much larger than those of the WT-β-RC and the Zn-β-RC (14% and 10%, respectively; Table 3). In our experiments, the degree of  $P^+$  remaining in the WT-β-RC at the end of the measurement, as indicated by the difference in bleaching amplitude at 70 nm between 45 ps post excitation and the end of the measurement window, was less than reported previously by Kirmaier et al. because our experiments used a longer measurement window (5.5 ns vs. 2.7 ns), which revealed additional amounts of ground state recombination [32]. Note that there is no evidence for any yield loss on the picosecond time scale in either the WT or mutant RCs. This can be seen by the invariance between the blue (few picoseconds) and green (hundreds of picoseconds) EADS near 830 nm in Fig. 6.



**Fig. 6.** Evolution associated difference spectra (EADS) of the  $Q_x$  and  $Q_y$  transitions of the RCs measured in this study. The solid black line is a zero absorbance reference line. Each RC was fit to a fast < 10 ps component, a hundreds of ps component, and a nanosecond or non-decaying component. Component lifetimes are representative of best fits to the transient absorption spectrum datasets using a least squares fitting approach. Deviations represent the overlap of 5% variations in the  $\chi^2$  (goodness of fit) values between independent fittings of the  $Q_x$  and  $Q_y$  datasets.

At this wavelength, there is essentially no contribution in the picosecond spectrum from stimulated emission, and so one can monitor the decay of ground state bleaching independently.

## 4. Discussion

### 4.1. Functional analysis of Zn- and Zn-β-RCs

We have exploited the Zn-BChl-containing RCs to evaluate the effects of changing the H<sub>A</sub> cofactor coordination state on electron transfer in the RC. This comparison would be impossible in the (Mg<sup>2+</sup>-) BChl-containing RC because the H<sub>A</sub> pocket contains BPhe when a fifth ligand is absent [20]. Evidently the Zn<sup>2+</sup> ion is more tightly bound to the bacteriochlorin macrocycle than Mg<sup>2+</sup>, and so Zn-BChl is resistant to whatever mechanism is responsible for the dechelation of BChl to yield BPhe [12,20,35].

The kinetic consequences of substituting Zn-BChl into the RC in an otherwise WT background were previously studied, and it was found that the Zn-RC functions similarly to the WT RC [9]. However, our EADS data now indicate a slower rate of electron transfer from H<sub>A</sub> to Q<sub>A</sub> for the Zn-RC. The EADS with a hundreds of ps time constant for the Zn-RC fitted best to ~290 ps, compared to ~210 ps in the WT RC. The difference between our experimental values and those previously obtained by Lin et al. likely stems from our use of a broadband detection setup, in contrast to a single wavelength detector [9]. Furthermore, the final charge separation yields of the Zn-RC were also lower, with evidence of some P<sup>+</sup>I<sub>A</sub><sup>-</sup> recombination in the non-decaying EADS. In this respect, the Zn-RC is more similar to the WT-β-RC and Zn-β-RC because the presence of ground state recombination on this time scale suggests the rate of recombination from P<sup>+</sup>I<sub>A</sub><sup>-</sup> is increased. The rate of recombination in the Zn-RC was thus calculated based on the observed lifetime of the P<sup>+</sup>I<sub>A</sub><sup>-</sup> state and the amount of P<sup>+</sup> remaining at the end of the time window of the experiment according to the formula:

$$\tau_{\text{rec}} = \tau_{\text{obs}} / ((100 - \phi) / 100) \quad (1)$$

where  $\tau_{\text{rec}}$  is the rate of recombination,  $\tau_{\text{obs}}$  is the P<sup>+</sup>I<sub>A</sub><sup>-</sup> lifetime, and  $\phi$  is the final yield of P<sup>+</sup> [32]. For the Zn-RC, this yielded a value of 1.1 ns, a value closer to that of the WT-β-RC and Zn-β-RC (651 ps and 711 ps, respectively) than the 20 ns value obtained for the WT-RC from which Q<sub>A</sub> had been removed (Table 3) [36].

The relatively long-lived hundreds of ps EADS obtained for the WT-β-RC from our datasets compared to previously published P<sup>+</sup>I<sub>A</sub><sup>-</sup> lifetimes suggests a degree of heterogeneity in the P<sup>+</sup>I<sub>A</sub><sup>-</sup> decay [32,37]. When our datasets were truncated to a shorter time window (2 ns), the P<sup>+</sup>H<sub>A</sub><sup>-</sup> decay rates were found to be closer to the previously reported rates (data not shown). The slower rates observed in the WT-β-RC and the Zn-β-RC likely reflect more complex kinetic heterogeneity than can be described by four exponential components, reflecting an average of forward electron transfer as well as ground state recombination from RCs in the P<sup>+</sup>I<sub>A</sub><sup>-</sup> state. A previous analysis of the kinetics of the WT-β-RC in which electron transfer to Q<sub>A</sub> was blocked demonstrated that the recombination rate from P<sup>+</sup>I<sub>A</sub><sup>-</sup> was approximately 1 ns (compared to ~20 ns in the WT-RC) [32]. Furthermore, our calculated charge recombination rates in the WT-β-RC, and Zn-β-RC suggest that charge recombination occurs much faster than in the WT-RC. This difference likely contributes to the relatively shallow fitting minima obtained in the WT-β-RC and Zn-β-RC for this EADS (compare the deviations between the WT-RC/Zn-RC and WT-β-RC/Zn-β-RC in the green spectra of Fig. 6). Because the Zn-RC also has a more rapid rate of charge recombination relative to the WT-RC, we suggest that the presence of the metal center at H<sub>A</sub> in the Zn-RC results in an H<sub>A</sub><sup>-</sup> anion that is closer in energy to B<sub>A</sub><sup>-</sup> compared to the WT-RC, as previous mutagenic studies on RCs have found a strong correlation between the rate of P<sup>+</sup>H<sub>A</sub><sup>-</sup> recombination and the similarity of B<sub>A</sub><sup>-</sup> and H<sub>A</sub><sup>-</sup> energies [38]. Clearly, this energetic proximity of H<sub>A</sub><sup>-</sup> to B<sub>A</sub><sup>-</sup> is lesser in the Zn-RC compared

to the WT-β-RC and Zn-β-RC, because the P<sup>+</sup>H<sub>A</sub><sup>-</sup> recombination rate of the Zn-RC is relatively slow.

In addition to the rate of recombination, it was also possible to extract information on the rates of forward electron transfer from I<sub>A</sub><sup>-</sup> to Q<sub>A</sub>, by taking into account the lifetime of the hundreds of ps EADS, as well as the final P<sup>+</sup>Q<sub>A</sub><sup>-</sup> yield of each RC. The rate of electron transfer was calculated according to the equation:

$$\tau_{\text{QT}} = \tau_{\text{obs}} / (\phi / 100) \quad (2)$$

where  $\tau_{\text{QT}}$  is the rate of electron transfer,  $\tau_{\text{obs}}$  is the P<sup>+</sup>I<sub>A</sub><sup>-</sup> lifetime, and  $\phi$  is the final yield of P<sup>+</sup> [32]. The values of  $\tau_{\text{QT}}$  are 0.2 ns for the WT-RC, 4 ns (WT-β-RC), 0.4 ns (Zn-RC), and 6.4 ns (Zn-β-RC). From these data, it appears that the presence of a fifth coordinate at the H<sub>A</sub> metal center is associated with a much reduced rate of forward electron transfer from I<sub>A</sub><sup>-</sup> to Q<sub>A</sub>.

The greater similarity of Zn-RC kinetic behavior to that of the WT-RC, than to the WT-β-RC and Zn-β-RC, evidently is because the H<sub>A</sub> cofactor is a 4-coordinate Zn-BChl. This indicates that the five-coordinate center present in the latter two RCs is the key player in the observed low yields of charge separation, and that the presence or absence of a metal in the H<sub>A</sub> bacteriochlorin is of less consequence. One explanation for the observed results stems from the Q<sub>x</sub> transition energies of the H<sub>A</sub> cofactors among these RCs. Previous studies on metal-substituted bacteriochlorins *in vitro* have reported a linear correlation between Q<sub>x</sub> transition energies, redox potential, and electronegativity (Pauling values) of the bacteriochlorin metal center [15,16]. We suggest that the 4-coordinate Zn-BChl at H<sub>A</sub> functions similarly to a BPhe molecule because the Zn-BChl transition energy, redox potential, and central ion electronegativity are closer to that of BPhe (as indicated by the Q<sub>x</sub> absorption peak) than a 5-coordinate BChl or Zn-BChl. In the case of central ion electronegativity, although the Pauling value for Zn<sup>2+</sup> (~1.6) is closer to Mg<sup>2+</sup> (~1.3) than H (~2), this value is similar to both and lies between Mg<sup>2+</sup> and H. Additionally, an increase in the coordination state of the metal center of a bacteriochlorin has been associated with a lower electronegativity value for the metal [16]. Thus, a 5-coordinate Zn-BChl should be less electronegative than both the 4-coordinate counterpart and BPhe.

In addition to the observed kinetic differences upon entering the P<sup>+</sup>I<sub>A</sub><sup>-</sup> state, the presence of a metal at H<sub>A</sub> imbues different spectroscopic properties on the I<sub>A</sub><sup>-</sup> anion. The transient absorption spectra of the P<sup>+</sup>I<sub>A</sub><sup>-</sup> state for Zn-BChl-containing RCs are more similar to each other (and the WT-β-RC) than the WT-RC in the 650–730 nm region. As noted in the Results section, the WT-RC has an absorption maximum at higher energy in this region compared to the other 3 RCs. Thus, the red-shifted peak in these three RCs appears to be a correlated with the presence of a metal at H<sub>A</sub>, regardless of coordination state or the type of metal present.

Interestingly, the RC of the photosynthetic bacterium *A. rubrum* contains Zn-BChl in place of BChl, but this organism contains BPhe at H<sub>A</sub>. In this organism, Zn-BChl is presumably synthesized by Mg-containing BChl synthesis intermediates [39]. After the final step of BChl synthesis, it is thought that the low environmental pH promotes pheophytinization BChl, at which point a Zn<sup>2+</sup> atom incorporates into the chlorin macrocycle by an unknown mechanism [39]. This theoretically provides a pool of both Zn-BChl and BPhe with which the apo-RC can selectively incorporate. This mechanism of Zn-BChl biosynthesis is different from that of the *R. sphaeroides bchD* mutant, because this mutant utilizes Zn-containing synthesis intermediates. Although electron transfer has been studied only to the P<sup>+</sup>H<sub>A</sub><sup>-</sup> state in *A. rubrum* [40], we speculate that in this RC the rate of electron transfer to Q<sub>A</sub> is like that in the WT-RC. Based on our results, it appears that natural selection would favor the presence of BPhe at H<sub>A</sub> in evolution because it results in a much slower decay of P<sup>+</sup>I<sub>A</sub><sup>-</sup> via charge recombination.

## 4.2. Consequences of the absence of Mg-BChl

In the ground state absorption spectra of the Zn-RC, as well as in previous low temperature measurements, a single peak was present in the  $Q_x$  region near 560 nm, originally attributed to overlapping absorption of 4-coordinate Zn-BChls in the  $H_{A,B}$  pockets [9,10]. However, as reported by Neupane et al. and confirmed here, the Zn- $\beta$ -RC has little or no absorption around 560 nm, whereas it would be thought that only the component due to the  $H_A$  Zn-BChl would be shifted as a result of the (M) L214H mutation, as is seen when this mutation is introduced into the WT RC (Fig. 1) [10]. Neupane et al. offered three possible explanations for the absence of a defined  $H_B$  peak in the  $Q_x$  region of the Zn- $\beta$  RC [10]. Our crystal structures show that, relative to the WT-RC, the Zn-BChl in the  $H_B$  pocket of the Zn-RC is bound at less than full occupancy or is highly disordered, and perhaps the occupancy is even lower or disorder greater in the Zn- $\beta$ -RC (Table 2). These conclusions are based on weaker density in omit difference maps of the  $H_B$  Zn-BChl, the large B-factors associated with the  $Zn^{2+}$  in the  $H_B$  pocket, and the low anomalous signal intensities for that  $Zn^{2+}$  atom (compare the values for  $H_B$  to those of the other cofactors in Table 2). Interestingly, the average B-factor of the  $H_B$  BPhe from the highest resolution RC structure currently available in the PDB database is greater than all of the other cofactors, perhaps indicating weaker binding to this pocket in the WT RC as well. As mentioned in the Results section, the B-factor for the protein component surrounding the  $H_B$  Zn-BChl is comparable to the average of the structure, consistent with the idea that the RC is a relatively rigid protein, and does not rearrange its folding due to changes in cofactor composition or even when a cavity is created by the lack of a cofactor, as was observed in pigment exclusion experiments on the  $H_B$  BPhe [41,42]. The larger  $Q_x$  spectral bandwidth of the  $H_B$  cofactor compared to  $H_A$  in is also supportive of a greater  $H_B$  site heterogeneity (Fig. 1). Additional evidence for relatively weak binding of BPhe in the WT RC  $H_B$  pocket comes from experiments involving substitution of BPhe with plant pheophytin (as well as a variety of other pigments). These experiments demonstrated that substitution is preferential in the  $H_B$  over the  $H_A$  site, consistent with a weaker binding of BPhe in the  $H_B$  pocket [43,44]. In vitro, the binding pockets for  $H_{A,B}$  are selective for pheophytins, as incubation of RCs in a molar excess of metal-containing bacteriochlorins (the standard procedure for pigment exchange) resulted in substitution of pigments at the  $B_{A,B}$  binding pockets instead of  $H_{A,B}$  [41,45,46]. The greater hydrophobicity of BPhe over BChl and Zn-BChl, as indicated by the longer retention time for BPhe in reversed-phase HPLC using a C18 column [35], may be part of the reason why binding of metal-containing bacteriochlorins to the  $H_{A,B}$  pockets is weak, although it is not clear why the  $H_B$  Zn-BChl may be lost more readily or less efficiently inserted into the RC than the  $H_A$  Zn-BChl. Alternatively, it is possible that the Zn-RCs contain partially occupied  $H_B$  sites due to a lack of a sufficient quantity of available Zn-BChls in the cell. One important consequence of the *bchD* mutation yielding Zn-BChl-producing strains of *R. sphaeroides* is the severely impaired ability to synthesize the pigment [35,47]. Therefore the inability of the organism to grow photoheterotrophically is evidently because of a paucity in the cellular content of light-harvesting and RC complexes [31].

Unlike the  $H_B$  cofactors of the Zn-RC and Zn- $\beta$ -RC, the  $H_A$  cofactors appear less disordered, suggesting there is no issue with the occupancy of the pocket. From a structural perspective, the  $H_A$  cofactors appear to assume a similar configuration in the two Zn-BChl-containing RCs (Fig. 2), with a possible macrocycle displacement of 0.6 Å between the two RCs. We are hesitant to make a firm conclusion on this geometric displacement between the cofactors because of the poor resolution of the data obtained for these two RCs; however the small degree of displacement is consistent with the interpretation that the differences in the functional characteristics of the two Zn-BChl-containing RCs stem from changes in  $H_A$  cofactor coordination state.

## 4.3. Biological implications of unnatural cofactor incorporation

Our time-resolved data demonstrate that in all 4 of these RCs  $P^+I_A^-$  is achieved rapidly (within a few ps), and that differences between RCs become apparent after this state is attained. Decay kinetics of the  $P^+H_A^-$  transient at 675 nm indicate that this state is relatively long-lived in the Zn- $\beta$ -RC, as well as in the WT- $\beta$ -RC, in agreement with previous results [11,27,32]. Furthermore, our global analysis suggests a similar degree of special pair charge recombination associated with this long-lived  $P^+H_A^-$  state in the Zn- $\beta$ -RC, much like in the WT- $\beta$ -RC. The charge separation yields are consistent with this interpretation (i.e., the WT- $\beta$ -RC and the Zn- $\beta$ -RC yields were 14% and 10%  $P^+Q_A^-$ , respectively, whereas the WT-RC and Zn-RC yields were 100% and 74%  $P^+Q_A^-$ , respectively). These results suggest that RCs have a degree of functional tolerance for different cofactors, and that the protein component is a key instrument in the design of an efficient charge separation pathway. The Zn-RC studied here is a case in point, as it contains a cofactor arrangement supportive of a faster rate of ground state recombination from the  $P^+I_A^-$  state, but is nevertheless capable of attaining a relatively high efficiency of charge separation. This is likely because the protein scaffold of the Zn-RC has remained unchanged from its WT-RC counterpart, granting this RC the same favorable cofactor geometries and tunneling pathways available to the WT-RC. From an evolutionary perspective, this could be advantageous to photosynthetic organisms during periods of environmental change, such as a change in the wavelengths of light available, for it would allow for the continued utilization of existing protein machinery when the production of a new type of (bacterio)chlorin would provide a selective advantage. In such a case, the new pigment could be incorporated into a protein “template” with suitable geometries, which would grant the organism some degree of charge separation capability, enabling photosynthetic growth. Mutations in protein side chains may later provide further energetic optimization, resulting in a high-efficiency RC that assembles with a new pigment. For this to happen, however, the photosynthetic organism must be able to survive for some time with an inefficient RC to allow beneficial mutations to accumulate. Previous results show that LH2-lacking *R. sphaeroides* strains containing RCs with small-nonpolar side chains at  $H_A$  retain their ability to grow at low photon fluxes [20], despite the relatively low efficiencies of these RCs (compared to the WT-RC) [22], indicating that purple photoheterotrophs are indeed quite capable of surviving with lower efficiency RCs.

## 5. Concluding remarks

The ability to study 4- and 5-coordinate metal-containing pigments at  $H_A$  provides insights into the use of (B)Phe vs. (B)Chls as the pigment of choice at this position in type II RCs. Firstly, the coordination state of the central metal of BChl studied here, and perhaps other metal-chelating tetrapyrroles, affects the electron transfer properties. This effect may be due to a decrease in energy and redox potential resulting from an increase in coordination number, as observed in protein-free BChls [15,16]. Secondly, the presence of a metal in the bacteriochlorin ring, irrespective of coordination state, appears to give rise to a longer lived  $P^+H_A^-$  state, however this change is minor in comparison to the introduction of a -coordinate center. Thirdly, the presence of a metal at  $H_A$  is associated with a faster rate of recombination to the ground state in the nanosecond time domain.

## Acknowledgements

This work was funded by NSF grants MCB-0642260 and MCB-1157788 at ASU. JTB and MEPM thank NSERC Canada for funding through the Discovery Grants system; RGS thanks NSERC for a postgraduate fellowship; AH thanks CIHR for a Canadian Graduate Fellowship. We thank Anson Chan for preparing Fig. 2. The spectrophotometer and cryostat were funded by a Canada Foundation for Innovation

Grant to AGM. Portions of this research were carried out at the Stanford Synchrotron Radiation Lightsource, a Directorate of SLAC National Accelerator Laboratory and an Office of Science User Facility operated for the U.S. Department of Energy Office of Science by Stanford University. The SSRL Structural Molecular Biology Program is supported by the DOE Office of Biological and Environmental Research, and by the National Institutes of Health, National Institute of General Medical Sciences (including P41GM103393) and the National Center for Research Resources (P41RR001209). The contents of this publication are solely the responsibility of the authors and do not necessarily represent the official views of NIGMS, NCRR or NIH.

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